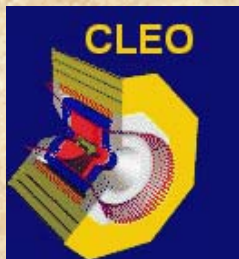


Semileptonic Charm Decays and QCD

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ISMD Sonoma
July 27, 2004



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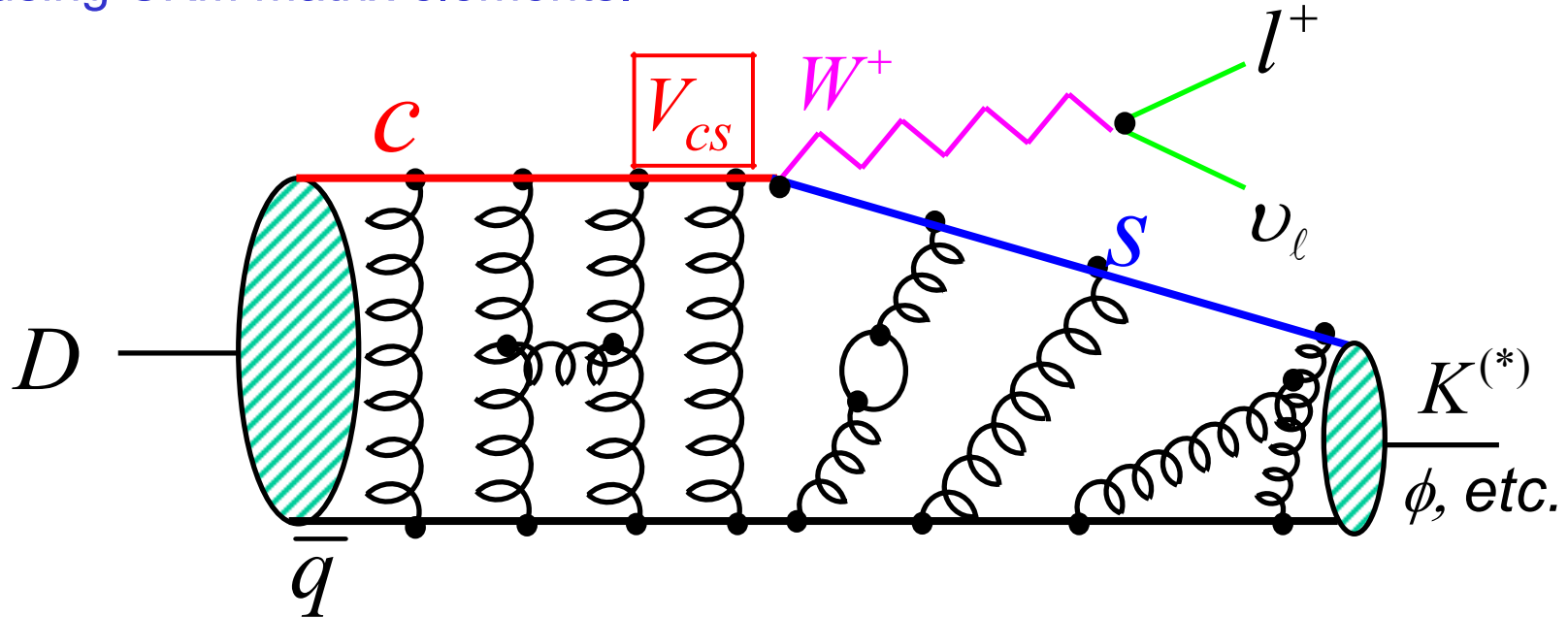
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Content

- Part I: Theories of Charm Semileptonic decays
- Part II: q^2 dependence in Pseudo-scalar $l \nu$ decays.
- Part III: Vector $l \nu$ decays.
 - $D^+ \rightarrow K^{*0} \mu \nu$ analysis (not so simple!)
 - Form factors of $D_s \rightarrow \phi \mu \nu$
- Part V: Future of Semileptonic decays.

I: Charm semileptonic decay as tests of LQCD

The decay rates are computed from first principles (Feynman diagrams) using CKM matrix elements.



The hadronic complications are contained in the form factors, which can be calculated via non-perturbative Lattice QCD, HQET or quark models.

Charm SL decays provide a high quality lattice calibration, which is crucial in reducing systematic errors in the Unitarity Triangle. The techniques validated by charm decays can be applied to beauty decays.

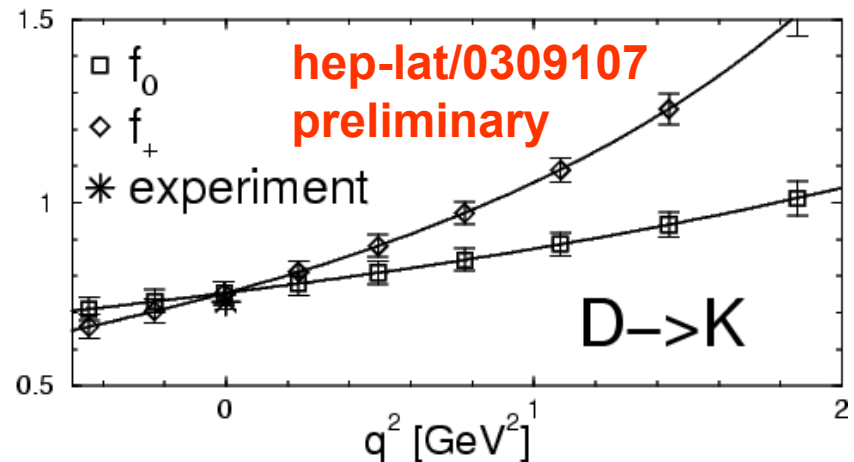
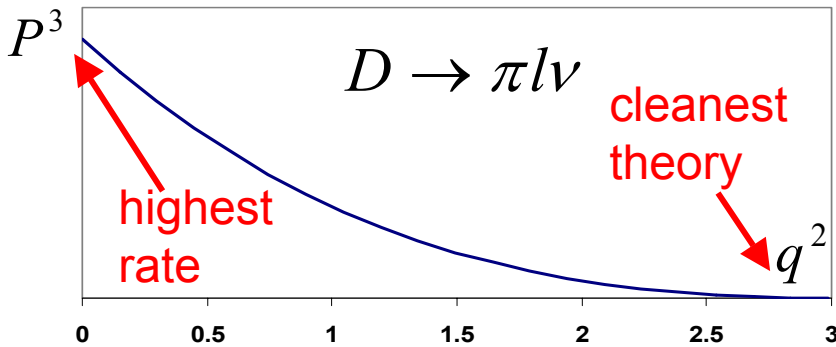
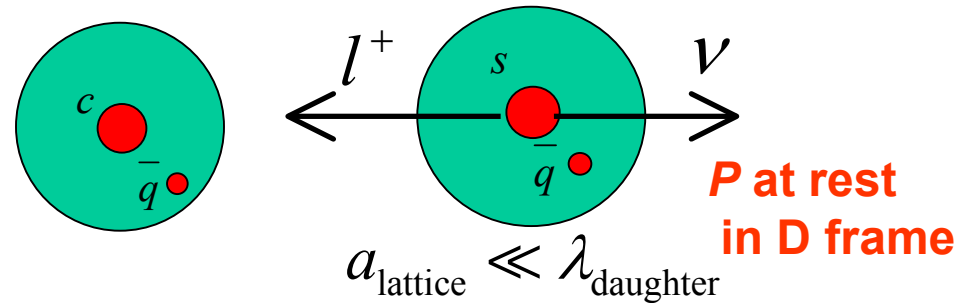
II: Pseudoscalar / ν decays

Simple kinematics
 → Easy to extract
 form factors.

$$\frac{d\Gamma(D \rightarrow P\ell\nu)}{dq^2} = \frac{G_F^2 |V_{cq}|^2 P_P^3}{24\pi^3} \left\{ |f_+(q^2)|^2 + O(m_\ell^2) \right\}$$

But a major disconnection exists between experiment and theory. In the past, theories worked best where experiments worked worst.

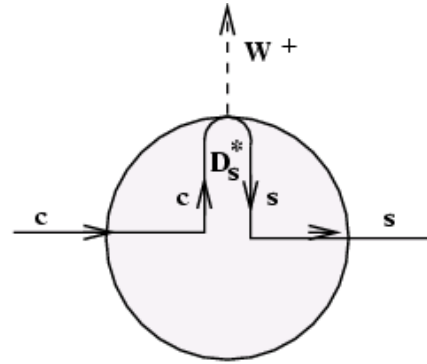
$q^2 = q_{\max}^2$ is the easiest point for LQCD calculation.



The lattice community is actively fixing the situation and calculating f_+ as a function of q^2 .

Comparing Pole versus ISGW forms in $D \rightarrow \pi l \nu$

Until quite recently one required a specific parameterized form to bridge the gap between a theory and an experiment, since neither an experiment nor a theory had clean $f_+(q^2)$ information.



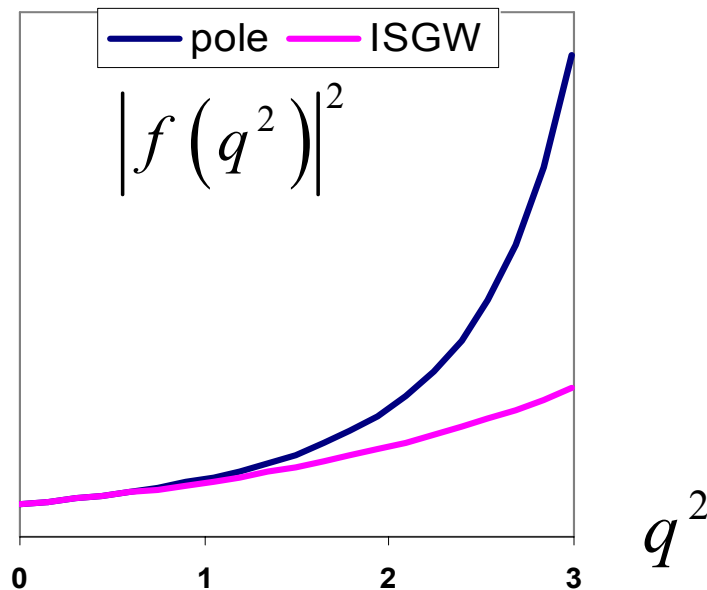
$f_+(q^2)$ parameterization

pole $f_+ \propto \frac{1}{q^2 - m_{\text{pole}}^2}$

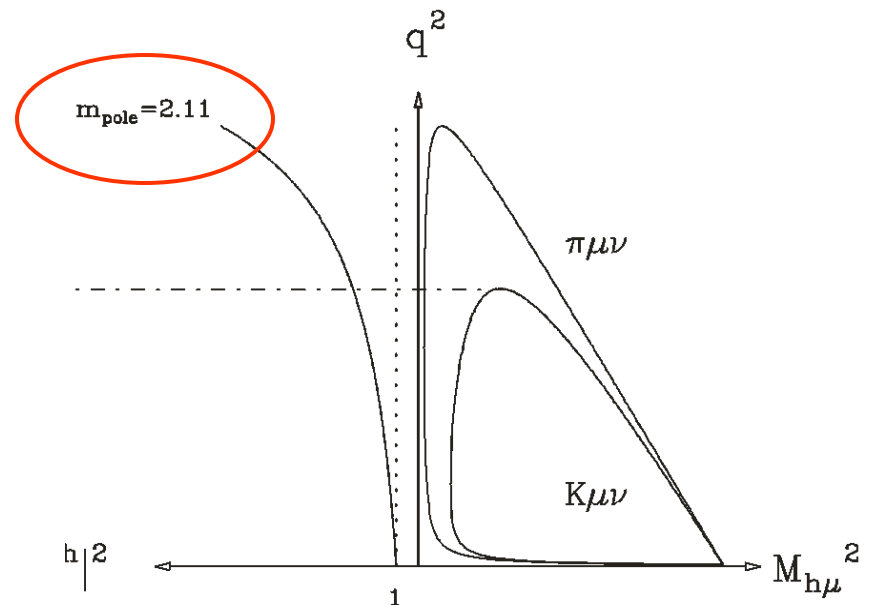
ISGW1 $f_+ \propto \exp(\alpha q^2)$

ISGW2 Updated one.

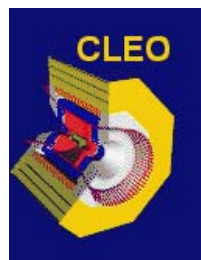
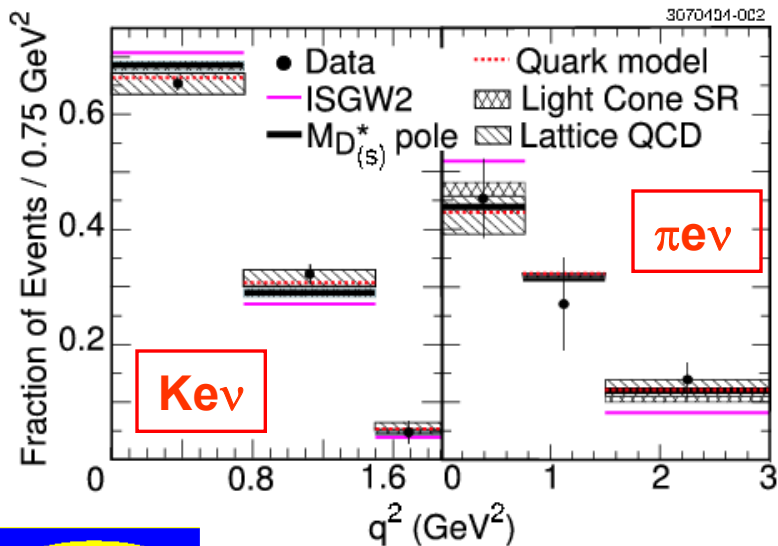
The difference between these forms can be quite dramatic in $\pi \mu \nu$ decays.



Especially since $\pi \mu \nu$ decay gets quite close to the D^* pole.



Brand new q^2 information in $D \rightarrow \pi l \nu / K l \nu$



Preliminary Cleo 2004
 $\pi \nu$ pole mass is

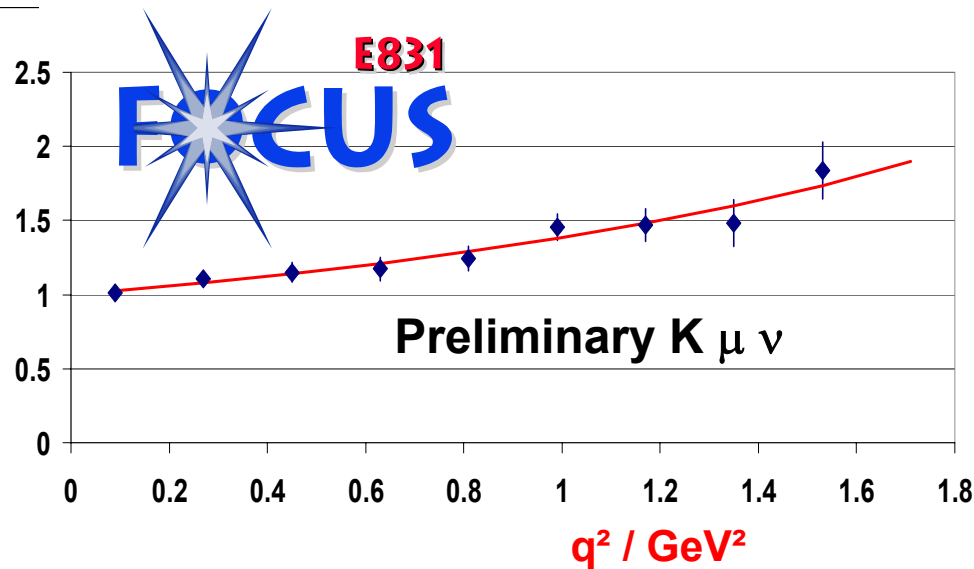
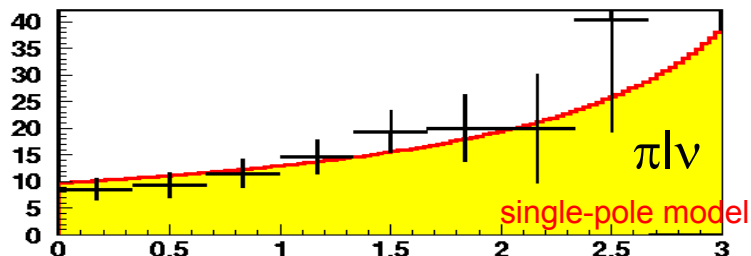
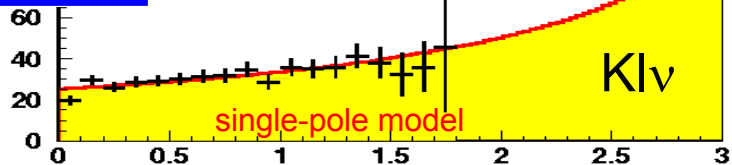
$$1.86^{+0.10+0.07}_{-0.06-0.03} \text{ GeV}$$

It disfavors ISGW2 form by $\sim 4.2\sigma$

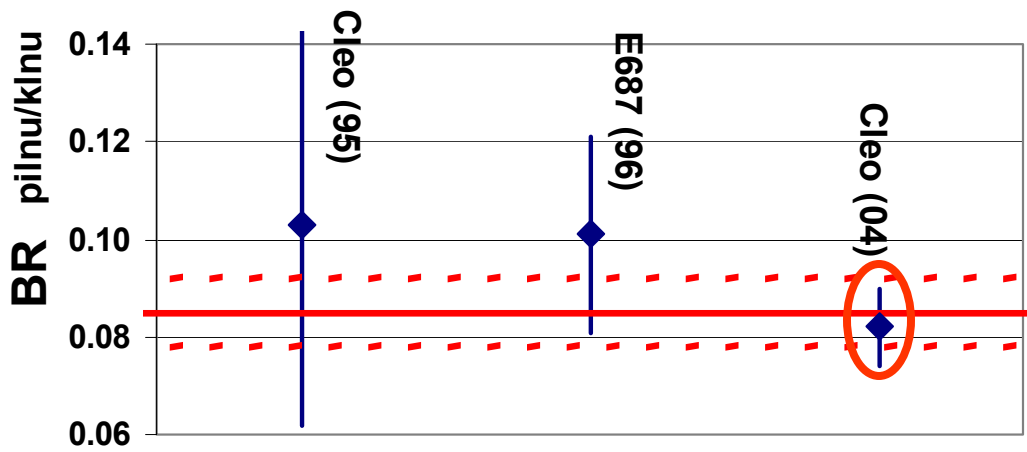
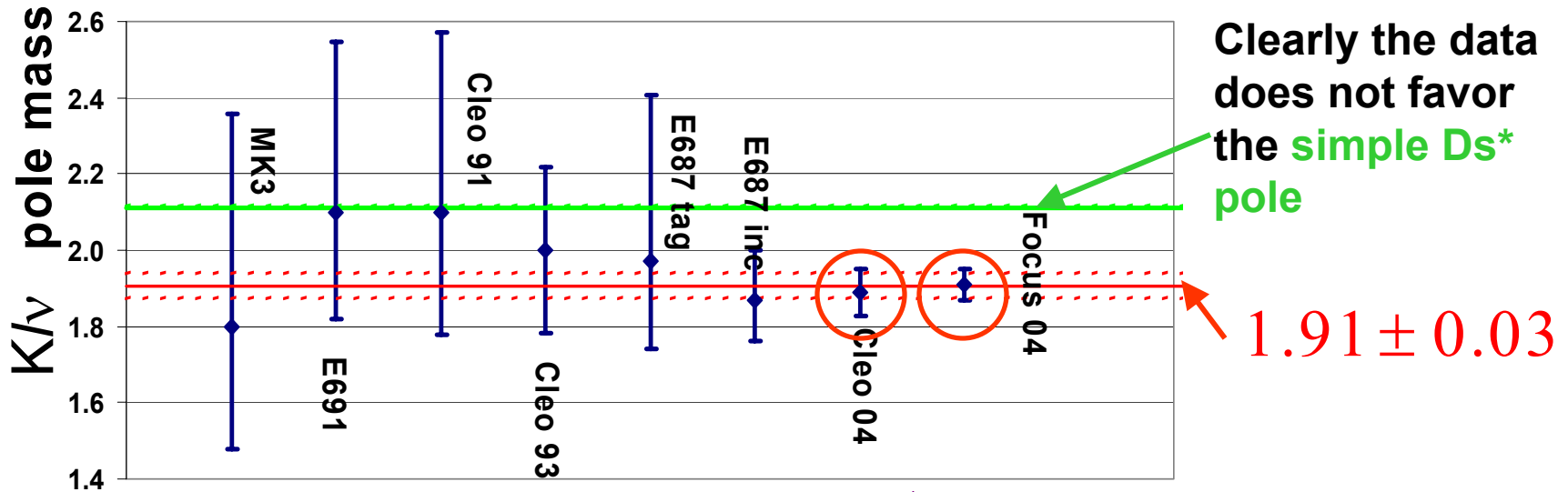


form factor $f_+(q^2)$

Based on 820 events



Summary of $D \rightarrow \pi l \nu / K l \nu$ Results



A big advance in precision!

$$\frac{d\Gamma(D \rightarrow Pl\nu)}{dq^2} = \frac{G_F^2 |V_{cq}|^2 P_P^3}{24\pi^3} |f_+(q^2)|^2$$

$$\frac{\Gamma(\pi e \nu)}{\Gamma(K e \nu)} = 0.082 \pm .006 \pm 0.005 \text{ CLEO}$$

$$\frac{|f_+^\pi(0)|}{|f_+^K(0)|} = 0.86 \pm 0.07 \pm 0.05 \pm 0.01$$

Consistent w/ SU(3) breaking

III: D → vector μ ν decays

Two amplitudes
get summed over W
polarization using
D-matrices

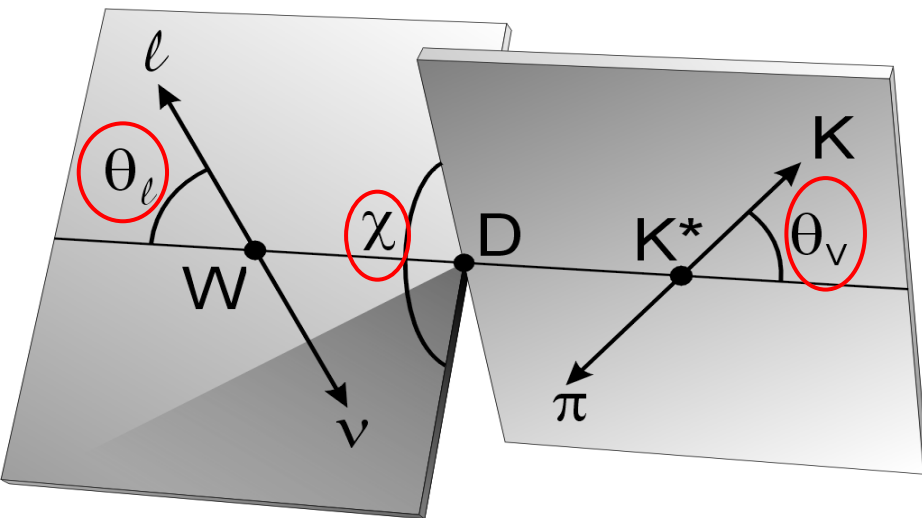
$$|A|^2 = \frac{1}{8} (q^2 - m_l^2)$$

right-handed μ^+

$$\left\{ \begin{array}{l} (1 + \cos \theta_l) \sin \theta_V e^{i\chi} H_+ \\ -(1 - \cos \theta_l) \sin \theta_V e^{-i\chi} H_- \\ -2 \sin \theta_l \cos \theta_V H_0 \end{array} \right\}^2 + \frac{m_\mu^2}{q^2} \left\{ \begin{array}{l} \sin \theta_l \sin \theta_V e^{i\chi} H_+ \\ + \sin \theta_l \sin \theta_V e^{-i\chi} H_- \\ + 2 \cos \theta_l \cos \theta_V H_0 \\ + 2 \cos \theta_V H_t \end{array} \right\}^2$$

left-handed μ^+

$H_0(q^2)$, $H_+(q^2)$, $H_-(q^2)$ are helicity-basis form factors computable by LQCD



Helicity FF are
combinations
of one vector
and two axial
form factors.

$$A_i(q^2) = \frac{A_i(0)}{1 - q^2/M_A^2}$$

$$V(q^2) = \frac{V(0)}{1 - q^2/M_V^2}$$

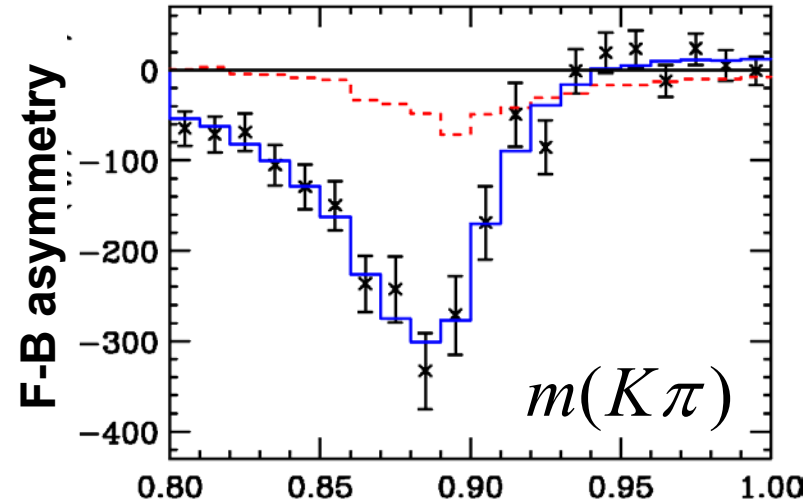
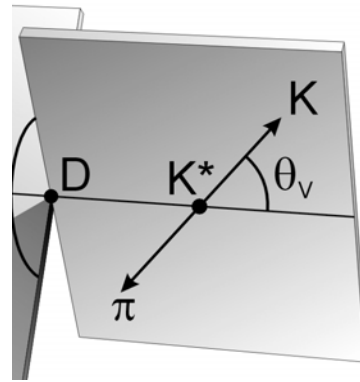
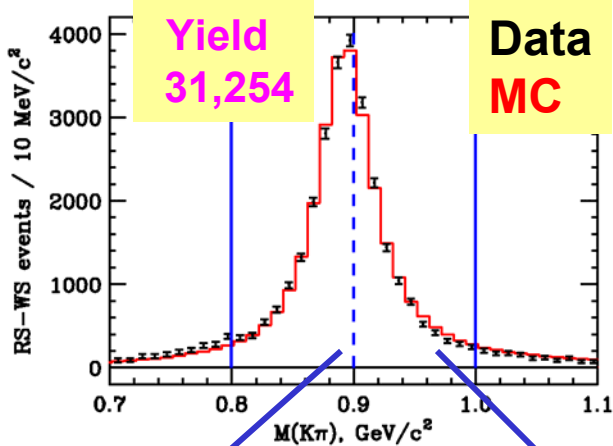
Two
observables
parameterize
the decay

$$r_V \equiv V(0)/A_1(0)$$

$$r_2 \equiv A_2(0)/A_1(0)$$

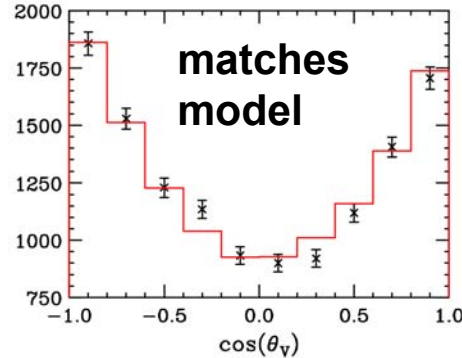
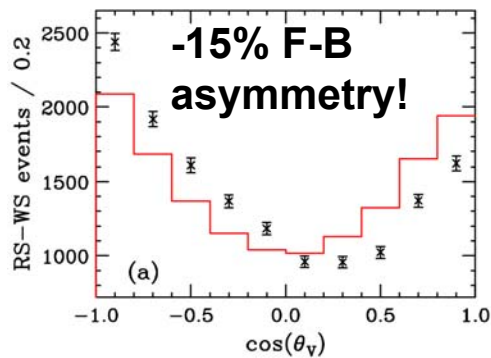
Interference in $D^+ \rightarrow K^* \mu \nu$

Focus “K*” signal



$0.8 < M(K\pi) < 0.9$ GeV/c²

$0.9 < M(K\pi) < 1.0$ GeV/c²



$K^* \mu \nu$ interferes with S-wave $K\pi$ and creates a forward-backward asymmetry in the K^* decay angle with a mass variation due to the varying BW phase.

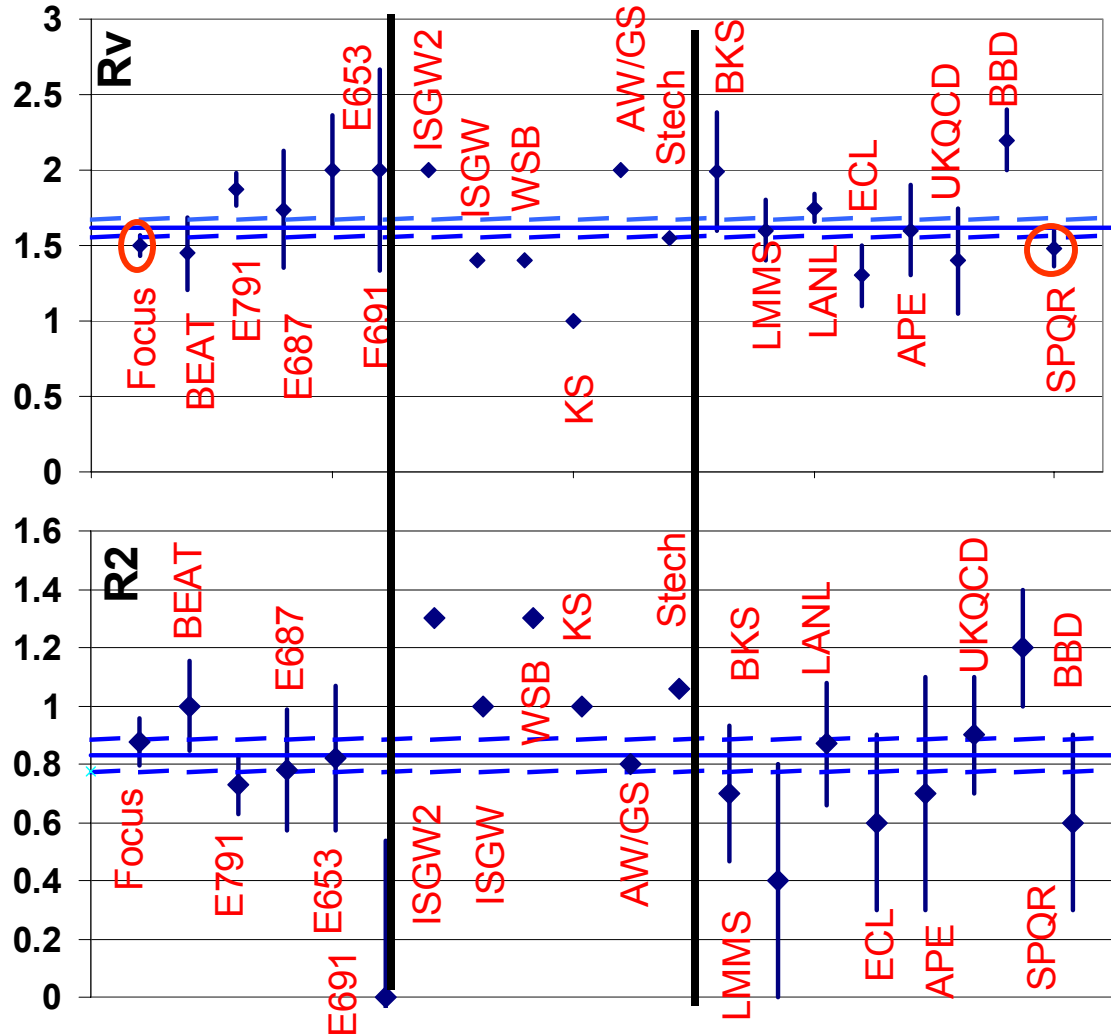
The S-wave amplitude is about 7% of the (H_0) K^* BW with a 45° relative phase

It's the same relative phase as the LASS (1988)



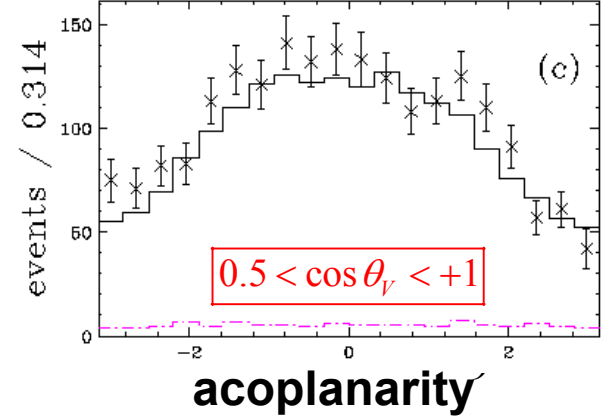
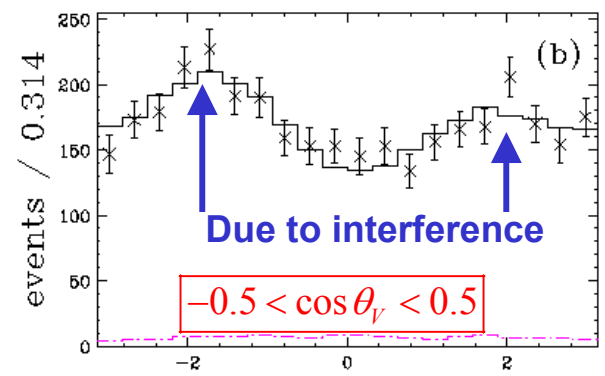
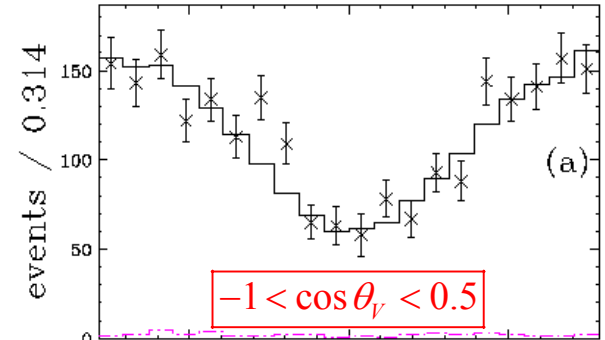
(2002)

$K^*_{\mu\nu}$ form factors



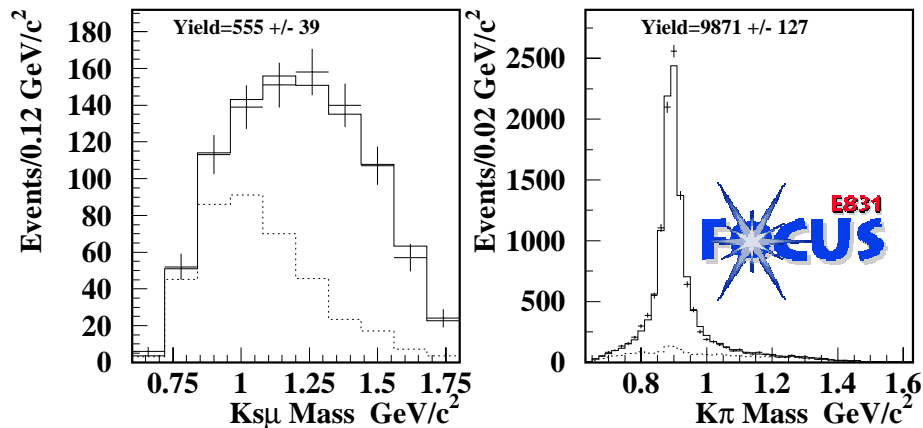
Results are getting very precise and unquenched calculations for incisive tests of the theory would be very desirable.

Precision tests of the model.

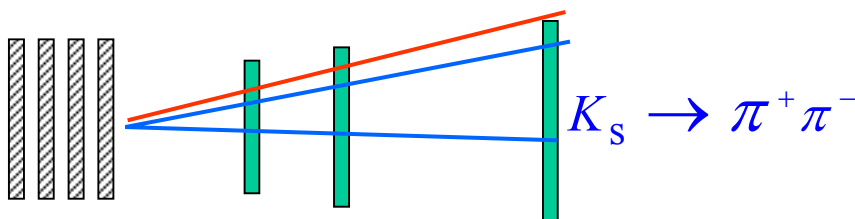


acoplanarity

Direct measurement of $\Gamma (D^+ \rightarrow K^* \mu \nu / K \mu \nu)$



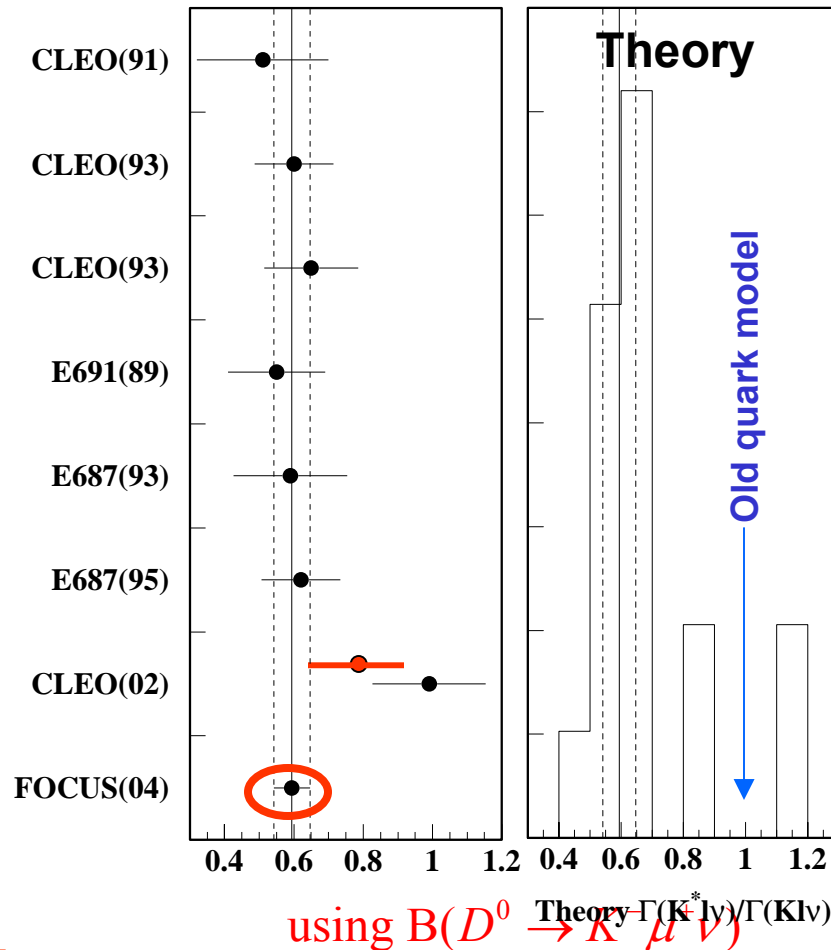
Use upstream K_s (~10%) so that both the numerator ($K \pi \mu \nu$) and denominator ($K_s \mu \nu$) leave 3 tracks in FOCUS μ -strip



$$\frac{D^+ \rightarrow \bar{K}^{*0} \mu^+ \nu}{D^+ \rightarrow \bar{K}^0 \mu^+ \nu} = 0.594 \pm 0.043 \pm 0.033$$

S-wave corrected

$$D^+ \rightarrow \bar{K}^{*0} l^+ \nu / \bar{K}^0 l^+ \nu$$



CLEO(02) partially reflects an inconsistency in $\Gamma (D^+ \rightarrow \bar{K}^0 e^+ \nu)$

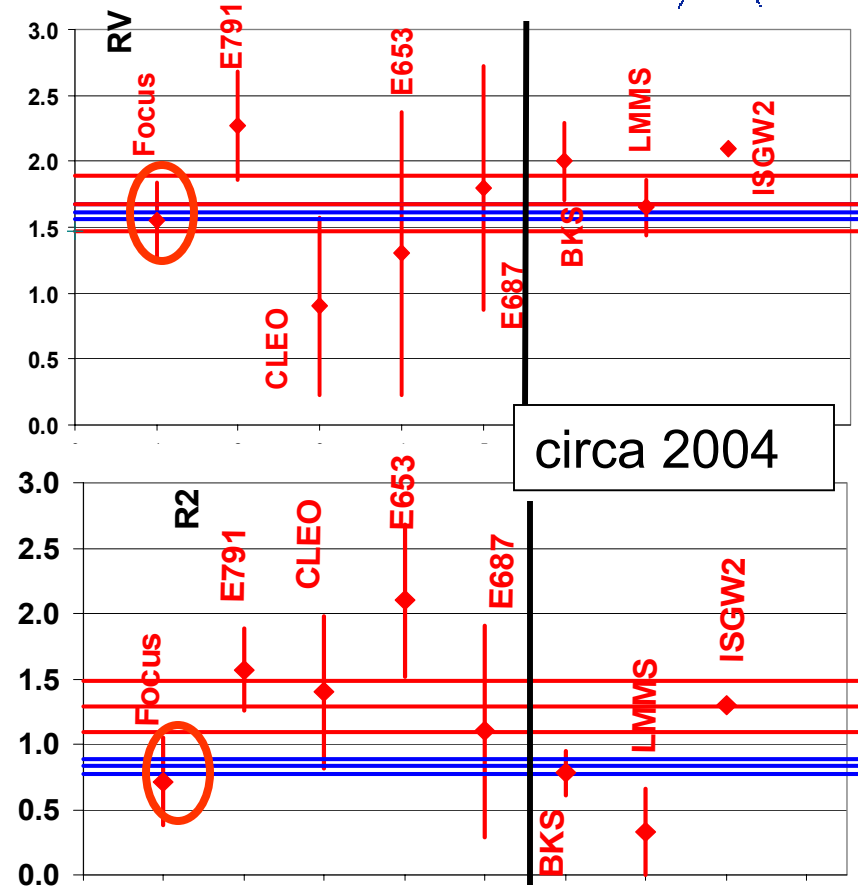
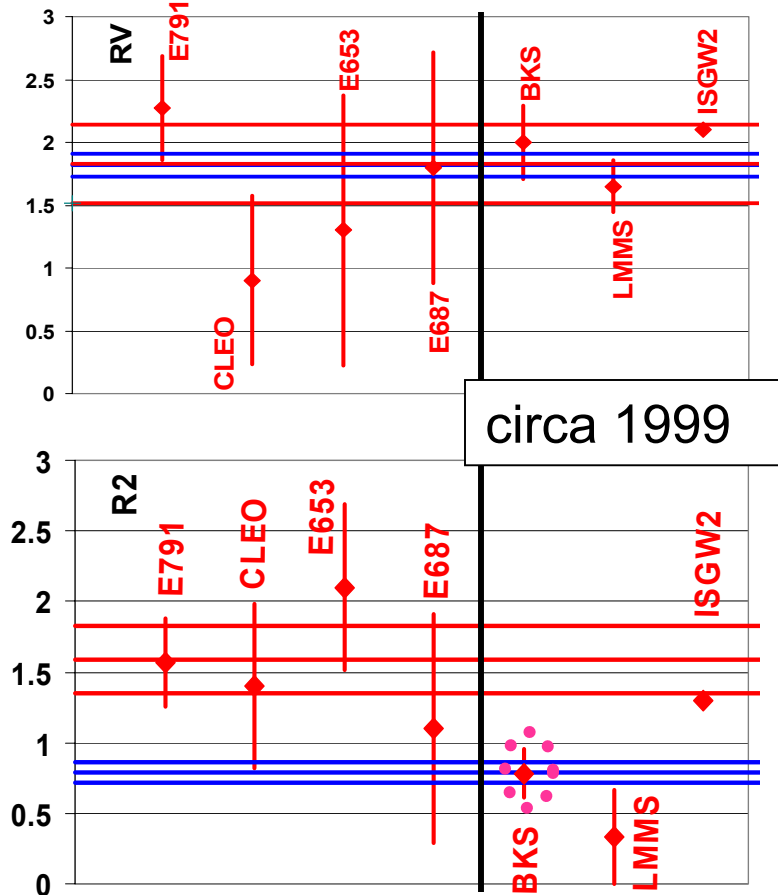
$$\Gamma_{\bar{K}^0 \mu^+ \nu}(D^+) - \Gamma_{K^- \mu^+ \nu}(D^0) = (11 \pm 11) / ns$$

$$\Gamma_{\bar{K}^0 e^+ \nu}(D^+) - \Gamma_{K^- e^+ \nu}(D^0) = (-25 \pm 9.7) / ns$$

The $D_s \rightarrow \phi \mu \nu$ form factor enigma



$D_s \rightarrow \phi \mu \nu$ versus $D^+ \rightarrow K^* l \nu$



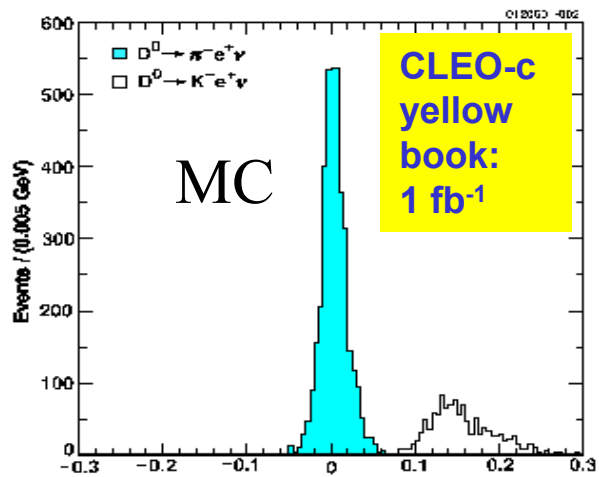
Theoretically, the $D_s \rightarrow \phi l \nu$ form factors should be within 10% of $D^+ \rightarrow K^* l \nu$. The r_V values were consistent, but r_2 for $D_s \rightarrow \phi l \nu$ was $2 \otimes$ higher than $D^+ \rightarrow K^* l \nu$.

But the (2004) FOCUS measurement obtained a consistent r_2 value as well!

The future of charm SL physics

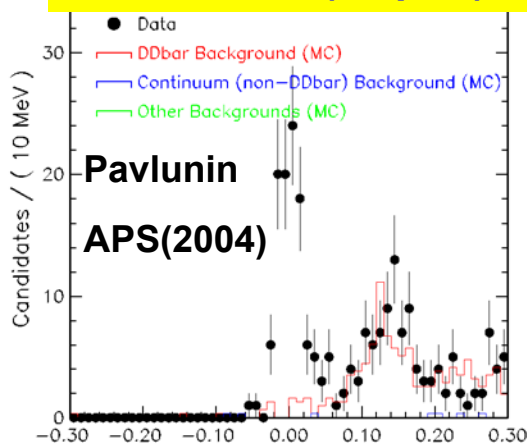
Cleo-c and Bes III: Run at $\Psi(3770)$ with high luminosity and modern detectors.

Precision neutrino closure in $D \rightarrow \pi e \nu$.



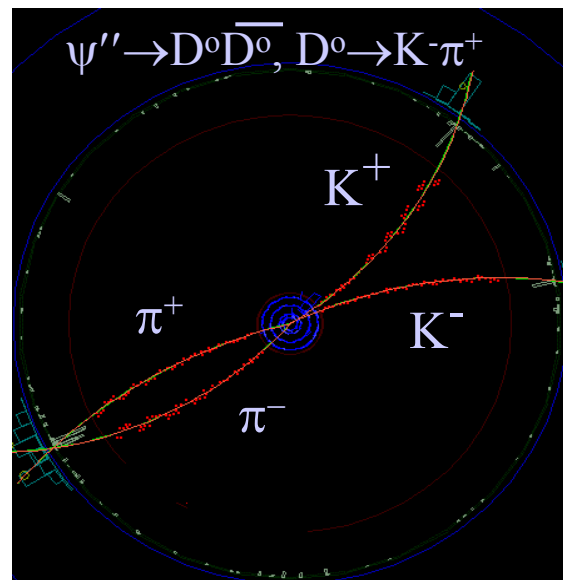
$$U = E_{\text{miss}} - P_{\text{miss}}$$

Prelim. data (60 pb⁻¹)

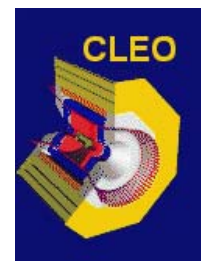
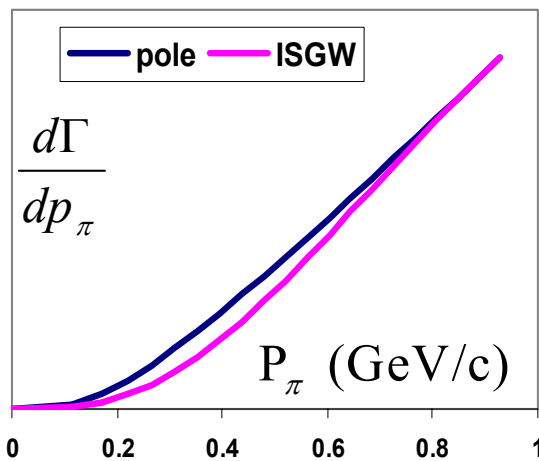
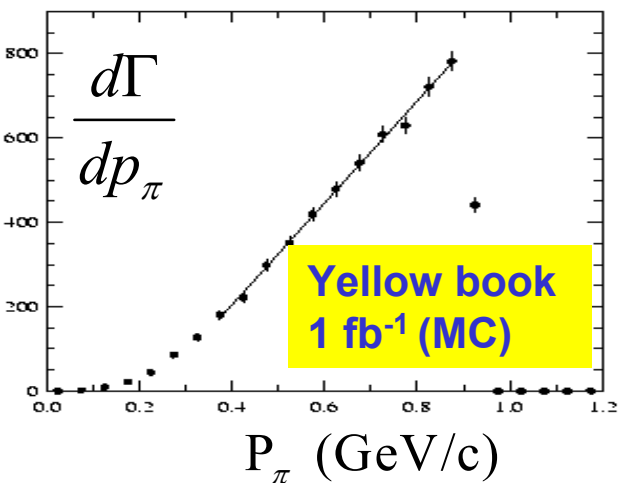


$$U = E_{\text{miss}} - P_{\text{miss}}$$

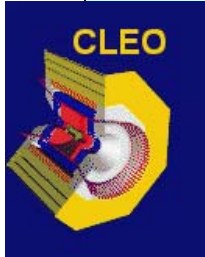
Extremely clean events!



The q^2 impasse afflicting SL data for the last 20 years shall be solved, finally.



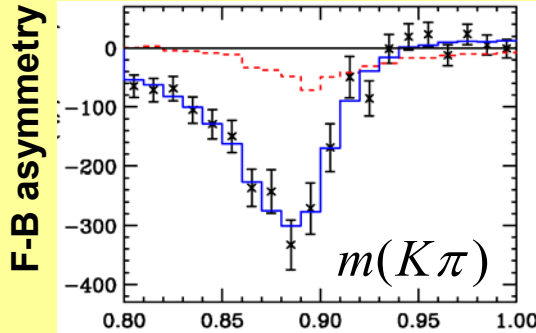
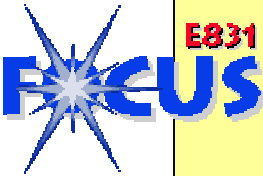
Summary



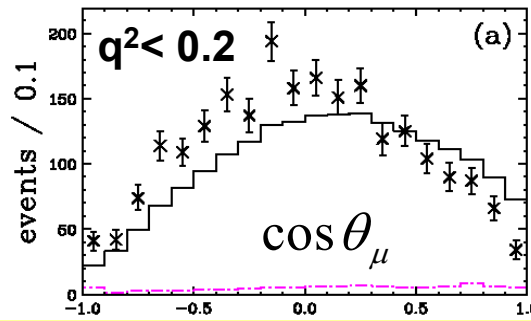
New CLEO 2004 $D \rightarrow \pi e \nu / K e \nu$ result

$$\frac{\Gamma(\pi e \nu)}{\Gamma(K e \nu)} = 0.082 \pm .006 \pm 0.005 \quad \left| \frac{f_+^\pi(0)}{f_+^K(0)} \right| = 0.86 \pm 0.07 \pm 0.05$$

$$m_{\text{pole}}^{D \rightarrow K} = 1.89_{-0.05}^{+0.05} {}_{-0.03}^{+0.04} \text{ GeV} < m(D_s^*)$$

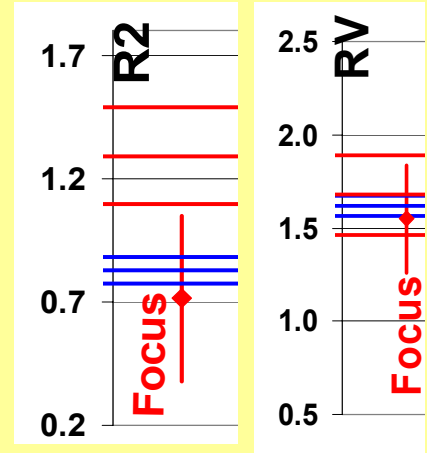
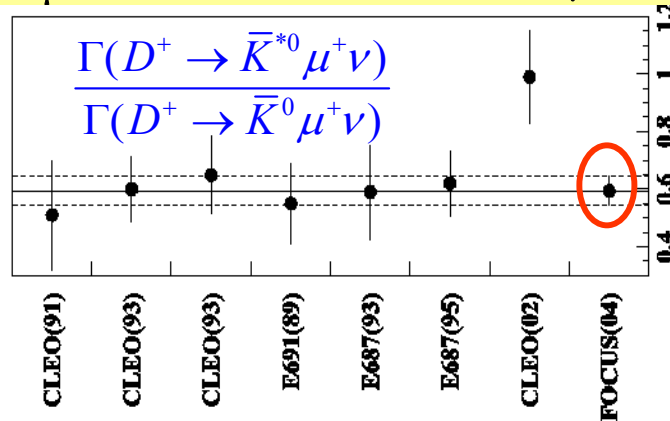


s-wave interference
in $D^+ \rightarrow K^* \mu \nu$



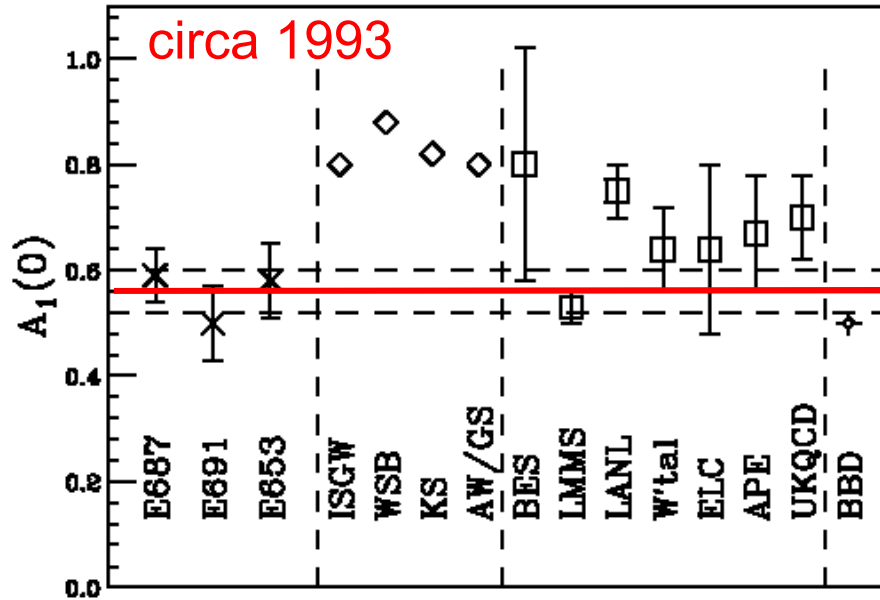
$V(q^2 \rightarrow 0)$ problem
 $D^+ \rightarrow K^* \mu \nu$

V/PS
ratio



Consistent FF for
 $D^+ \rightarrow K^* \mu \nu$ & $D_s^+ \rightarrow \phi \mu \nu$

$\Gamma (D^+ \rightarrow K^* \mu \nu / K \mu \nu)$ circa 1993



Some more tests of the $K^*\mu\nu$ model

A dramatic mismatch is seen at very low q^2 suggesting a $V(q^2 \rightarrow 0)$ problem



Generally the model tracks the data rather well...

Focus has a preliminary analysis of the K^{*0} line shape. $\Gamma(K^{*0})$ is seen as less than PDG by ~ 1.6 MeV.

